



Computational Fluid Dynamic- Two Fluid Model Study of Gas-Solid Heat Transfer in a Riser with Various Inclination Angles

K. Azizi, M. Keshavarz Moraveji*

Department of Chemical Engineering, Amirkabir University of Technology (Tehran Polytechnic), Tehran, Iran

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ABSTRACT

A two fluid model (TFM) was used to study gas-solid heat transfer in a riser with different inclination angles. A two dimensional pipe with 5.8 cm internal diameter and 5 meter length was chosen. Effect of bed angle and solid particles feed rate were studied on the heat transfer behavior of gas and solid particles. Obtained results from simulation are compared with the experimental data in the relevant literature. Heat transfer behavior of phases is different in an inclined pipe in comparison with vertical and horizontal pipes. It is found that higher air-solid Nusselt number, air temperature difference and particle temperature difference take place at the pipe with inclination angle equal to 45 degrees. Loading ratio enhancement increases gas temperature differences. At lower and higher loading ratio, particles temperature differences decreases and increases respectively with loading ratio enhancement.

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NOMENCLATURE

C_D	Drag coefficient	Greek Symbols	
C_p	Specific heat	α	Volume fraction
d_p	Solid diameter	β_{gs}	Gas-solid interphase exchange coefficient
e_{ss}	Coefficient of restitution of the particle	γ_{Θ_s}	Collision dissipation of energy
g	Gravitational constant	γ_{gs}^0	Heat-transfer coefficient
$\bar{G}_{k,g}$	Production of turbulent kinetic energy	ε	Turbulence dissipation rate
H_{gs}	Heat transfer between gas and solid	η_{gs}	Ratio between characteristic times
I	Turbulent intensity	Θ	Granular temperature
I_{2D}	Second invariant of the deviatoric stress	$\mu_{s,col}$	Collisional part of the shear viscosity
k	Turbulent kinetic energy tensor	$\mu_{s,fr}$	Shear viscosity
k'	Thermal conductivity	$\mu_{t,g}$	Turbulent (or eddy) viscosity
k_{Θ_s}	Diffusion coefficient for granular energy	$\mu_{s,kin}$	Kinematics viscosity
Nu_s	Nusselt number	v	Velocity
p	Pressure	$\Pi_{k,g}, \Pi_{\varepsilon,g}$	Influence of the dispersed phase on the continuous phase
Pr	Gas Prandtl number	$\bar{\tau}$	Reynolds stress tensor
Re	Reynolds number	Subscripts	
T	Time	g	Gas
V	Velocity	s	Solid
		t	Turbulence

*Corresponding Author's Email: Moraveji@aut.ac.ir (M. Keshavarz Moraveji)

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1. INTRODUCTION

Solids moving with a gas stream in a pipeline can be found in many industrial processes. These include chemical, pharmaceutical, food, mineral industries, gaseous and particulate pollutant transport in the atmosphere and many other applications [1]. In many of these applications, heat transfer between gas and solid particles in the mixture and to the surroundings has a significant impact on the performance of the system [2]. High surface area of particles contributes to high heat transfer rates [3]. As a consequence, there has been widespread interest in the effect of heat transfer on the design of these systems [2].

Three main modes of heat transfer, bed to surface, interparticle and gas to particle heat transfer arise in gas solid flows and are important to study [4]. Heat transfer studies have been done for gas-solid flows in the literatures. Shi et al. [5] studied wall to bed heat transfer in a CFB. They showed that solid concentrations and solid movements in the vicinity of the heat transfer surface have a significant influence on the wall to bed heat transfer coefficient. Rajan et al. [6] studied the effect of solids feed rates, air velocity and particle size on the heat transfer of gas-solid in the vertical pneumatic conveyor. Mansoori et al. [4] studied turbulent gas solid flows and showed experimentally that variation of Nusselt number can be explained by the variation of the thermal turbulence intensity. Despite a significant amount of research on fluidized bed reactors, there are considerable uncertainties on their behavior. With the increased computational capabilities, computational fluid dynamic (CFD) has become an important tool for understanding these complex phenomena that occur between the gas phase and the particles in fluidized bed [1]. Two different approaches have been developed for application of CFD to gas-solid flows; Eulerian-Lagrangian method and Eulerian-Eulerian method [7-9]. Simulation of these systems with the Eulerian models using different CFD codes has shown the suitability of the approach for modeling these systems [1]. Azizi et al. [10] used a two fluid model based on kinetic theory to study the heat transfer of gas solid flows in a vertical pneumatic conveyor and showed that heat transfer coefficient of gas –solid flow is greater than clean gas. Hamzehei and Rahimzadeh [11] used a multifluid Eulerian model incorporating the kinetic theory for solid phase to study the effect of particle size on heat transfer and hydrodynamics of a nonreactive gas-solid fluidized bed. Rajan et al. [12] simulated countercurrent gas-solid flow using two fluid model in a heat exchanger and compared their results with experimental data. They showed that heat transfer rates increase with increasing solid loading ratio and decreasing particle size. Behjat et al. [13] carried out a study on the gas-solid fluidized bed applying CFD

technique in order to investigate hydrodynamic and heat transfer phenomena. Gas-solid heat transfer during dilute phase pneumatic conveying was numerically investigated using an Eulerian-Eulerian approach by Rajan et al. [14]. Convective heat transfer between gas and particles and between gas and pipe wall, and conductive heat transfer between particles and pipe wall was modeled using distinct element method (DEM) by Li and Mason [2]. In their work, the gas phase was modeled as a continuum and the solid phase was modeled by DEM. Effect of collision-induced alterations of the flow dynamic on heat transfer in gas-solid flows in a vertical pipe have been presented by Chagras et al. [15].

In this study, the heat transfer of a two dimensional nonreactive gas-solid flow in a riser with different inclination angles was studied computationally. Computational study of riser with different inclination angles had been done for first time in this paper. A two fluid model (TFM) of multiphase incorporating the kinetic theory for solid particles with the $k - \epsilon$ turbulent model was applied to simulate the gas-solid flow. Momentum exchange coefficient was evaluated using Gidaspow drag functions. The results of simulation were compared with the experimental data of Mokhtarifar et al. [16] for model validation.

2. MODELING

The Eulerian-Lagrangian discrete element model (DEM) and the Eulerian-Eulerian two-fluid model (TFM) have been developed for CFD modeling of gas-solid flows. In the DEM, the fluid is treated as a continuum phase by solving the Navier-Stokes equations, while the dispersed phase is solved by tracking a large number of particles through the flow field. The dispersed phase can exchange momentum, mass and energy with the fluid phase.

In the TFM, both phases are treated as interpenetrating continuum phases [17]. In this method, the conservation equations for each of the two phases are derived. Therefore, two sets of coupled equations with similar mathematical structure should be solved. The success of this method depends on proper description of the forces between and acting on the phases.

In the present work, the governing equation of solid and gas phases are based on the Eulerian-Eulerian model. The governing equations of the gas-solid flow include the conservation of mass, momentum and energy [18, 19].

2. 1. Conservation of Mass

The continuity equation for gas and solid phases in the absence of interphase mass transfer are respectively given as:

$$\frac{\partial}{\partial t}(\alpha_g \rho_g) + \nabla \cdot (\alpha_g \rho_g \vec{v}_g) = 0 \quad (1)$$

$$\frac{\partial}{\partial t}(\alpha_s \rho_s) + \nabla \cdot (\alpha_s \rho_s \vec{v}_s) = 0 \quad (2)$$

where α_s and α_g are the volume fraction of solid and gas phase respectively. ρ is density and v is velocity.

2. 2. Conservation of Momentum The conservation of momentum for the gas and solid phases are described by:

$$\frac{\partial}{\partial t}(\alpha_g \rho_g \vec{v}_g) + \nabla \cdot (\alpha_g \rho_g \vec{v}_g \vec{v}_g) = -\alpha_g \nabla p_g + \nabla \cdot \bar{\tau}_g + \alpha_g \rho_g \vec{g} + \beta_{gs}(\vec{v}_s - \vec{v}_g) \quad (3)$$

$$\frac{\partial}{\partial t}(\alpha_s \rho_s \vec{v}_s) + \nabla \cdot (\alpha_s \rho_s \vec{v}_s \vec{v}_s) = -\alpha_s \nabla p_g - \nabla p_s + \nabla \cdot \bar{\tau}_s + \alpha_s \rho_s \vec{g} + \beta_{gs}(\vec{v}_g - \vec{v}_s) \quad (4)$$

Here, $\bar{\tau}$ is Reynolds stress tensor, g is the gravitational constant, and $(-\nabla p_s + \beta_{gs}(\vec{v}_g - \vec{v}_s))$ is an interaction force (drag and buoyancy) representing the momentum transfer between gas and solid phases. This term is a statement of Darcy's law where the reciprocal of β_{gs} is the permeability decided by the fluid viscosity [20]. Several drag models for gas-solid interphase exchange coefficient, β_{gs} , is reported in the literature [10]; the drag model of Syamlal-O'Brien, Gidaspow and Cao-Ahmedi. In this study, the drag model of Gidaspow was used.

The drag model of Gidaspow is combination of two drag model; Wen and Yu model for dilute flow and Ergun equation for dense flow [10]. When the volume fraction of gas, α_g , is greater than 0.8, β_{gs} is calculated with equations from the Wen and Yu model.

$$\beta_{gs} = \frac{3}{4} C_D \frac{\alpha_g \alpha_s \rho_g |\bar{v}_s - \bar{v}_g|}{d_s} \alpha_g^{-2.65} \quad (5)$$

The drag coefficient, C_D , is related to the Reynolds number by

$$C_D = \frac{24}{Re_s} (1 + 0.15(Re_s)^{0.687}), \quad Re_s < 1000 \quad (6)$$

$$C_D = 0.44, \quad Re_s \geq 1000 \quad (7)$$

$$Re_s = \frac{\alpha_g \rho_g (v_g - v_s) d_p}{\mu_g} \quad (8)$$

For dense flow ($\alpha_g \leq 0.8$), β_{gs} is calculated with Ergun equation.

$$\beta_{gs} = 150 \frac{\alpha_s^2 \mu_g}{\alpha_s d_s^2} + 1.75 \frac{\alpha_s \rho_g |\bar{v}_g - \bar{v}_s|}{d_s} \quad (9)$$

The constitutive equation for solid phase stresses is given by: [1]

$$\bar{\tau}_s = \alpha_s \mu_s [\nabla \vec{v}_s + (\nabla \vec{v}_s)^T] + \alpha_s \left(\lambda_s - \frac{2}{3} \mu_s \right) \nabla \vec{v}_s \bar{I} \quad (10)$$

The granular temperature is proportional to the specific kinetic energy of the random fluctuating component of the particle velocity. It should be mentioned that granular temperature is quite different from solid-phase thermodynamic temperature. The transport equation for the solid-phase granular temperature is given as: [1]

$$\frac{3}{2} \frac{\partial}{\partial t} (\alpha_s \rho_s \Theta_s) + \nabla \cdot (\alpha_s \rho_s \vec{v}_s \Theta_s) = (-p_s \bar{I} + \bar{\tau}_s) : \nabla \vec{v}_s + \nabla \cdot (k_{\Theta_s} \nabla \Theta_s) - \gamma_{\Theta_s} + \phi_{gs} \quad (11)$$

The first term of the right hand side of this equation is the generation of energy by the solid stress tensor; the second term is the diffusion flux of granular energy (k_{Θ_s} is the diffusion coefficient); the third term is the collisional dissipation of energy; and the fourth term is the energy exchange between the gas and solid. γ_{Θ_s} which represents the rate of energy dissipation due to inelastic particle collisions is given as: [1]

$$\gamma_{\Theta_s} = \frac{12(1-e_{ss}^2)g_{0,ss}}{d_s \sqrt{\pi}} \alpha_s^2 \rho_s \Theta_s^{3/2} \quad (12)$$

Random fluctuations cause the transfer of kinetic energy which is expressed as

$$\phi_{gs} = -3\beta_{gs}\Theta_s \quad (13)$$

The solid phase pressure, p_s , consists of a collisional and a kinetic part [21]:

$$p_s = \alpha_s \rho_s \Theta_s + 2\rho_s (1 + e_{ss}) \alpha_s^2 g_{0,ss} \Theta_s \quad (14)$$

where, e_{ss} is the restitution coefficient for particle collisions and $g_{0,ss}$ is the radial distribution.

$$g_{0,ss} = \left[1 - \left(\frac{\alpha_s}{\alpha_{s,max}} \right)^{1/3} \right]^{-1} \quad (15)$$

The solid phases shear viscosity consists of a collisional term:

$$\mu_{s,col} = \frac{4}{5} \alpha_s \rho_s g_{0,ss} (1 + e_{ss}) \left(\frac{\Theta_s}{\pi} \right)^{1/2} \quad (16)$$

A kinetic term:

$$\mu_{s,kin} = \frac{\alpha_s \rho_s d_s \sqrt{\Theta_s \pi}}{6(3-e_{ss})} \left[1 + \frac{2}{5} (1 + e_{ss}) (3e_{ss} - 1) \alpha_s g_{0,ss} \right] \quad (17)$$

And a friction term:

$$\mu_{s,fr} = \frac{p_s \sin \phi}{2\sqrt{I_{2D}}} \quad (18)$$

This parameter plays a significant role when the solid phase volume fraction gets close to the packing limit [1].

Hamzehei et al. [1] suggested the following equation which is solid bulk viscosity explaining the resistance of the granular particles to compression and expansion.

$$\lambda_s = \frac{4}{3} \alpha_s \rho_s d_s g_{0,ss} (1 + e_{ss}) \sqrt{\frac{\Theta_s}{\pi}} \quad (19)$$

k_{Θ_s} , the diffusion coefficient for the granular energy, is expressed by two the following equation [1].

$$k_{\Theta_s} = \frac{15\alpha_s\rho_s d_s \sqrt{\Theta_s \pi}}{4(41-33\eta)} \left[1 + \frac{12}{5}\eta^2(4\eta-3)\alpha_s g_{0,ss} + \frac{16}{15\pi}(41-33\eta)\eta\alpha_s g_{0,ss} \right] \quad (20)$$

where:

$$\eta = \frac{1}{2}(1 + e_{ss}) \quad (21)$$

2. 3. Internal Energy Conservation Neglecting radiation heat transfer, the internal energy balances for the gas and solid phase are given by:

$$\alpha_g \rho_g C_{pg} \left(\frac{\partial T_g}{\partial t} + v_g \cdot \nabla T_g \right) = -H_{gs} \quad (22)$$

$$\alpha_s \rho_s C_{ps} \left(\frac{\partial T_s}{\partial t} + v_s \cdot \nabla T_s \right) = \nabla \cdot \alpha_s K_s \nabla T_s + H_{gs} \quad (23)$$

The gas heat conductivity is negligible, so the heat diffusion term has been ignored. Solid granular conductivity is expressed by Ahmadi model [1].

$$K_s = 0.1306\rho_s d_s (1 + e_{ss}^2) \left(\frac{1}{g_{0,ss}} + 4.8\alpha_s + 12.1184g_{0,ss}\alpha_s^2 \right) \sqrt{\Theta_s} \quad (24)$$

The heat transfer between the gas and the solid is:

$$H_{gs} = -\gamma_{gs}^0 (T_s - T_g) \quad (25)$$

The heat transfer coefficient is related to the particle Nusselt number as follows.

$$\gamma_{gs}^0 = \frac{6k_g \epsilon_s Nu_s}{d_s^2} \quad (26)$$

k_g is the thermal conductivity of the gas phase. The Nusselt number is proposed by applying the Gunn correlation [13].

$$Nu_s = (7 - 10\alpha_g + 5\alpha_g^2)(1 + 0.7Re_s^{0.2}Pr^{1/3}) + (1.33 - 2.4\alpha_g + 1.2\alpha_g^2)Re_s^{0.7}Pr^{1/3} \quad (27)$$

where, Re_s is the relative solid Reynolds number and Pr is the gas Prandtl number.

2. 4. Turbulent Model For turbulent flows, viscosity in the momentum conservation equation is modified by a turbulence model which adds a turbulent viscosity to fluid of laminar flow. Viscosity of fluid is a property of fluid while the turbulent viscosity depends on flow conditions [10].

The $k - \epsilon$ model includes the turbulent kinetic energy (k) and the dissipation rate (ϵ) [22]. The Reynolds stress tensor for the gas phase is:

$$\bar{\tau}_g = -\frac{2}{3}(\rho_g k_g + \rho_g \mu_{t,g} \nabla \cdot \vec{v}_g) \bar{I} + \rho_g \mu_{t,g} [\nabla \vec{v}_s + (\nabla \vec{v}_s)^T] \quad (28)$$

The turbulent kinetic energy (k) and the dissipation rate (ϵ) are obtained from the following equations [10].

$$\frac{\partial}{\partial t} (\alpha_g \rho_g k_g) + \nabla \cdot (\alpha_g \rho_g k_g \vec{v}_g) = \nabla \cdot \left(\alpha_g \frac{\mu_{t,g}}{\sigma_k} \nabla k_g \right) + \alpha_g G_{k,g} - \alpha_g \rho_g \epsilon_g + \alpha_g \rho_g \Pi_{k,g} \quad (29)$$

$$\frac{\partial}{\partial t} (\alpha_g \rho_g \epsilon_g) + \nabla \cdot (\alpha_g \rho_g \epsilon_g \vec{v}_g) = \nabla \cdot \left(\alpha_g \frac{\mu_{t,g}}{\sigma_\epsilon} \nabla \epsilon_g \right) + \alpha_g \frac{\epsilon_g}{k_g} (C_{1\epsilon} C_{k,g} - C_{2\epsilon} \rho_g \epsilon_g) + \alpha_g \rho_g \Pi_{\epsilon,g} \quad (30)$$

where, $\Pi_{k,g}$ and $\Pi_{\epsilon,g}$ represent the influence of the solid phase (dispersed phase) on the gas phase (continuous phase). The constitutive equations and coefficients are given in literature [1].

2. 5. Initial and Boundary Condition The initial values of volume fraction and velocity of phases are specified for the entire computational domain. Initially, velocity of solid particle was set at zero. The temperature of both solid and gas phases was set at 300 K. Velocities and volume fraction of all phases were specified at the inlet of riser. A no slip wall condition at the pipe wall and constant pressure at outlet were used corresponding to momentum conservation equation. The boundary condition for the energy equation was set such that the wall is adiabatic. Table 1 shows the values of the parameters of the model used in the simulation.

3. RESULTS AND DISCUSSION

Simulation of the flow conditions is carried out using MFIX code. The 2D schematic of the pipe used in this simulation is shown in Figure 1. The computational domain used in this work is the same as the experimental work of Mokhtarifar et al. [16]. They used an inclined pipeline with length and height of $D=0.058\text{m}$ and $H=5\text{m}$, respectively. Obtained results from simulation are compared with the experimental data in the relevant literature [16] to verify the availability of presented model. The domain was divided into $20*500$ computation cells. The simulation was continued for 5 s and time step was set to 0.0001 s.

3. 1. Effects of Loading Ratio and Pipe Angle on Air and Particle Temperature Variations of axial air and particles temperature difference in different loading ratio and different pipe angles for $253 \mu\text{m}$ particle size, with air velocity of 18.5 m/s, are shown in Figure 2.

It is observed that gas temperature difference increase as loading ratio increase. Also, it can be seen from the figure that temperature difference of gas increase as the pipe angle increase until it reaches to 45 degrees and then temperature difference of gas decrease by changing the angle of pipe from 45 degrees to 90 degrees. The trends of results obtained from simulation and obtained by Mokhtarifar et al. [16] are in agreement with each other.

TABLE 1. Values of Model Parameters Used in the Simulations and Experiments

Property	Value
Particle density	1500 kg/m ³
Particle specific heat	800 J/kg-k
Particle thermal conductivity	0.80 W/m-k
Particle diameter	253 μm
Gas density	1.225 kg/m ³
Restitution coefficient	0.99
Superficial gas velocity	18.5 m/s
Bed width	5.8 cm
Bed height	5 m
Solid feed rate	9, 15, 18, 22, 25 g/s
Pipe angle	0, 30, 45, 60, 90 degrees
Initial temperature	298 k
Air inlet temperature	440 k
Solid inlet temperature	300 k

With increasing inclination angles, the difference between simulation and results obtained by Mokhtarifar et al. [16] increased. However, this difference decrease as the pipe angle increased from 45 to 90. Also, it can be seen from the figure that temperature difference in horizontal and vertical configuration has less dependency on loading ratio.

It is showed that maximum pressure drop in a pipe takes place at an angle lower than 90 degrees [16]. Pressure drop in a pipe causes more interaction between phases and higher heat transfer between phases and consequently higher particle temperature.

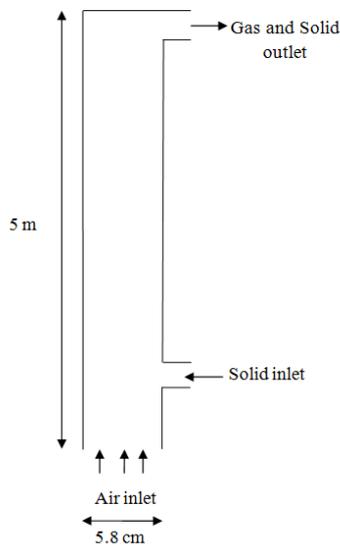


Figure 1. The 2D schematic of the modeling domain

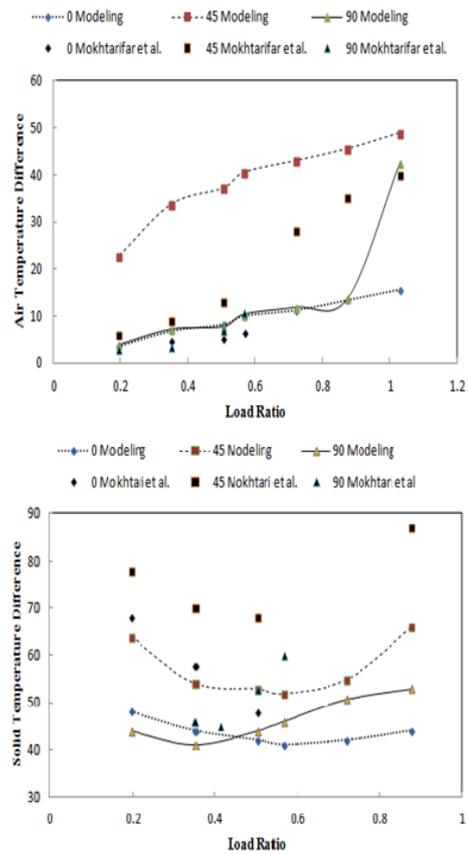


Figure 2. Variation of air and particle temperature difference versus loading ratio at different pipe angles

Contour plot of solid volume fraction, pressure and temperature of gas phase are shown in Figures 3-5. Also, gas pressure and temperature, before and after feeder are tabulated in Tables 2-3. In Figures 3-5, it could be seen that with changing the inclination angle the solid volume fraction and pressure drop will change. These changes will enhance heat transfer as it could be clearly seen from the Figures 6 (a-c). These results are in agreement with results presented in the literatures [16].

Similarly, temperature difference of particles increase as the pipe angle increase until it reaches to 45 degrees and then this difference decrease by changing the angle of pipe from 45 degrees to 90 degrees.

TABLE 2. Solid volume fraction before and after feeder in different inclination angles

Inclination Angle	Solid Volume Fraction Before Feeder	Solid Volume Fraction After Feeder
Horizontal	0.00003	0.00064
45 Degrees	0.00004	0.01656
Vertical	0.00003	0.00069

TABLE 3. Pressure drop before and after feeder in different inclination angles

Inclination Angle	Before Feeder	After Feeder
Horizontal	0.11528E+07	0.11526E+07
45 Degrees	0.11560E+07	0.11549E+07
Vertical	0.11531E+07	0.11528E+07

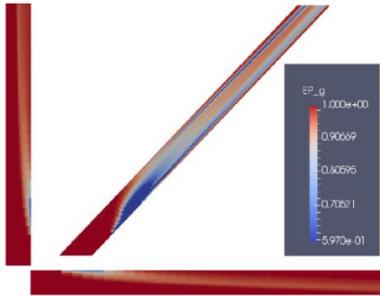


Figure 3. Contour plot of axial solid volume fraction at different pipe angles

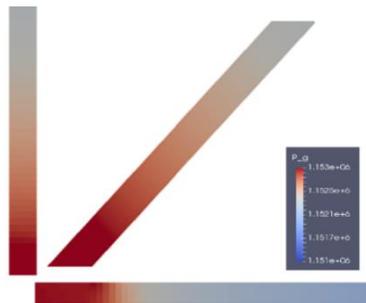


Figure 4. Contour plot of axial air pressure drop at different pipe angles

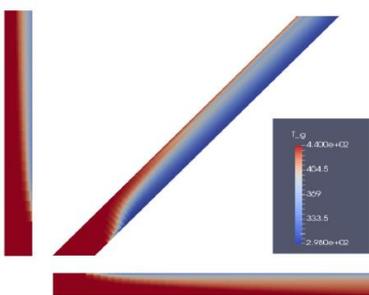


Figure 5. Contour plot of axial air temperature at different pipe angles

However, changing of particle temperature difference with loading ratio is different with what was obtained for air. It is observed that particles temperature decreases with increase of loading ratio at lower loading ratios, while it increases at higher loading ratios. It is showed that particle temperature is respectively proportional to air temperature difference and loading ratio directly and inversely [16].

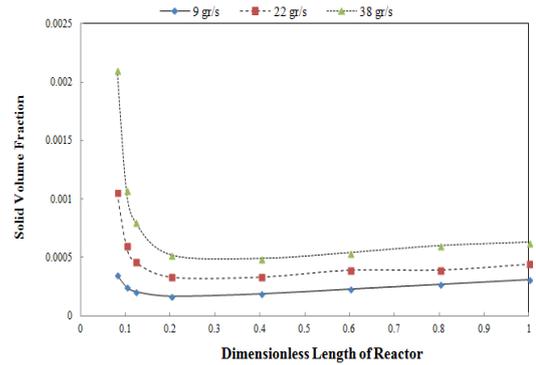


Figure 6a. Axial solid volume fraction in a horizontal pipe

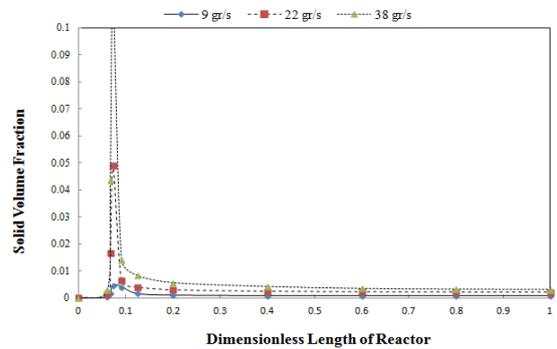


Figure 6b. Axial solid volume fraction in a pipe with inclination angle of 45 degrees

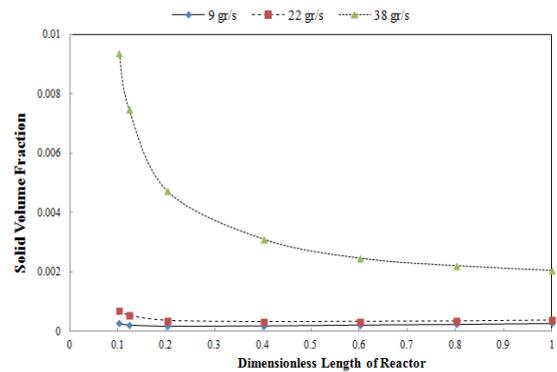


Figure 6c. Axial solid volume fraction in a vertical pipe

Thermal behavior of particles changes at a specific air velocity and a specific location. The pipe is divided into two parts in this location which are known as the dense phase and the dilute phase at the bottom and at the top of pipe. The height of dense phase starts to increase by increasing solid feed rate [16]. Obtained results from simulation are presented in Figure 6. In this figure, y axis is showing solid volume fraction and x axis is showing dimensionless length of the pipe. The results are obtained for three different inclination angles and different solid feed rates. It could be concluded from the figures that in a specific height of the pipe the

solid volume fraction in a higher particle feed rate is greater than that in a lower particle feed rate.

3. 2. Effects of Loading Ratio and Pipe Angle on Air-Particle Nusselt Number

The effect of the loading ratio on the air-particle Nusselt number for different pipe angles is illustrated in Figure 7. Air-particle Nusselt number is changed slightly by Air-particle heat transfer coefficient if air heat conductivity is constant. Also, it is showed that at lower solids feed rates the heat transfer coefficient decreases with increase of loading ratio and at higher solids feed rates it increases with increase of loading ratio. Obtained results are in good agreement with results observed by Mokhtarifar et al. [16].

It can be seen from the Figure 7, Nusselt number increases as the pipe angle reaches to 45 and then it decrease by changing the angle of pipe from 45 degrees to 90 degrees.

As it is showed the Nusselt number is directly proportional to the heat transfer rate and inversely proportional to driving force and heat transfer area. It could be concluded that at the pipe angle of 45 degrees with a constant air velocity and particle feed rate, heat transfer rate has its maximum value and driving force is minimum [16]. Consequently, the Nusselt number would be maximized.

3. 3. Thermal Effectiveness of Air and Solid

Figures 8 (a-c) show thermal effectiveness of air at different loading ratio and different pipe angles. It is seen that thermal effectiveness of air increases as the loading ratio increase. Thermal effectiveness is directly proportional to air temperature difference. On the other hand, it is inversely proportional to inlet driving force. In this simulation, driving force which is defined as difference between air and articles initial temperature is constant. Therefore, it was expected that thermal effectiveness and air temperature difference have the same behavior as it is shown in Figure 2.

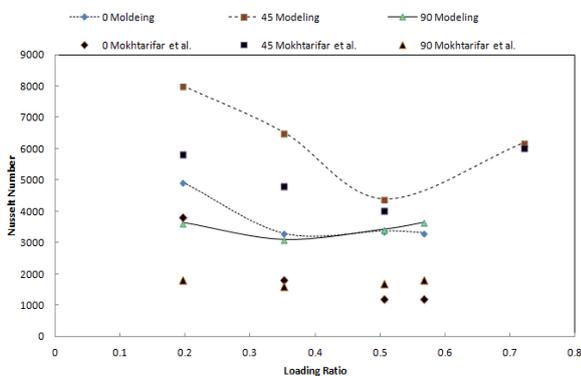


Figure 7. Variation of air - particle Nusselt number versus loading ratio at different pipe angles

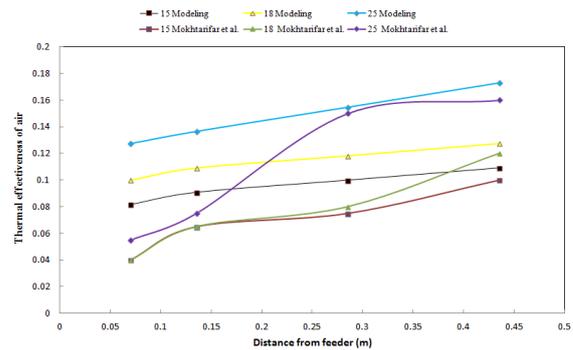


Figure 8a. Variation of thermal effectiveness of air versus distance from feeder for horizontal pipe at different particles loading ratios.

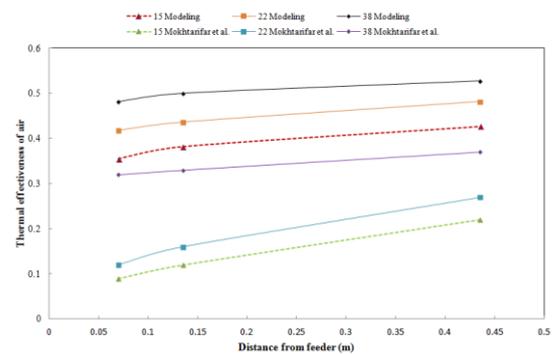


Figure 8b. Variation of thermal effectiveness of air versus distance from feeder for inclined pipe at different particles loading ratios.

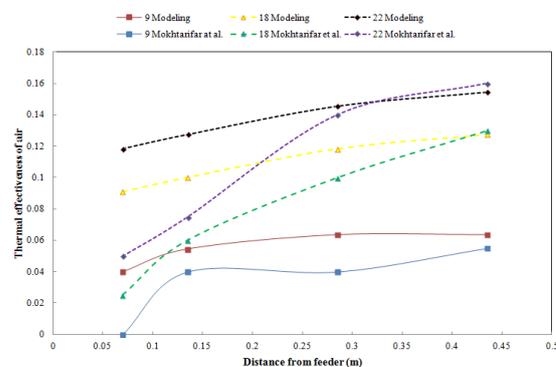


Figure 8c. Variation of thermal effectiveness of air versus distance from feeder for vertical pipe at different particles loading ratios.

The effect of feed rate on thermal effectiveness of solid for different pipe angles is illustrated in Figure 9. It is showed that solid thermal effectiveness is directly proportional to particle temperature difference and inversely proportional to driving force at inlet [16]. Since, inlet driving force is constant, it could be concluded that the behavior of thermal effectiveness of particles should be like the behavior of particle temperature at different loading ratio. This result can be clearly observed in Figure 9.

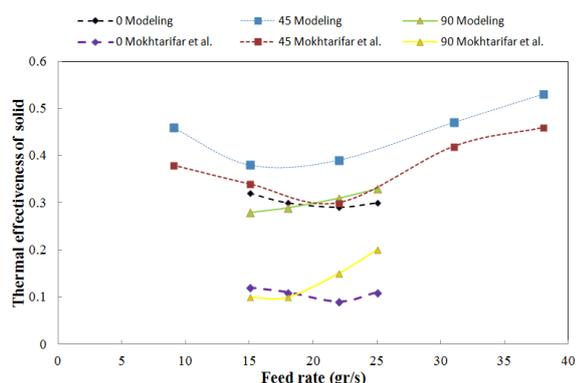


Figure 9. Variation of thermal effectiveness of solid versus solid loading ratio at different pipe angles.

4. CONCLUSION

A two fluid model (TFM) of multiphase flows based on the kinetic theory was used for numerical modeling of gas-solid flow and heat transfer in a fluidized bed with different inclination angles and particles loading ratios. Obtained results from simulation are compared with experimental data. It is observed that heat transfer behavior of phases is different in an inclined pipe in comparison with vertical and horizontal pipes. Higher Nusselt numbers take place at pipe angles closer to 45 degrees. Also, it is observed that gas temperature differences increase as loading ratio increases. On the other hand, particles temperature difference decreases with increase of loading ratio at lower loading ratios, while it increases at higher loading ratios.

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Computational Fluid Dynamic- Two Fluid Model Study of Gas-Solid Heat Transfer in a Riser with Various Inclination Angles

K. Azizi, M. Keshavarz Moraveji

Department of Chemical Engineering, Amirkabir University of Technology (Tehran Polytechnic), Tehran, Iran

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در این مقاله مدل دو سیالی به منظور بررسی انتقال حرارت بین دو فاز گاز و جامد در یک رایزر با زوایای مختلف نسبت به افق به کار برده شد. به این منظور، یک لوله با ارتفاع ۵ متر و قطر داخلی ۵/۸ سانتی متر بررسی شد. اثر زاویه بستر با افق و نرخ ورود ذرات جامد به بستر بر روی انتقال حرارت بین فاز گاز و ذرات جامد مطالعه شد. نتایج به دست آمده از شبیه سازی با نتایج تجربی ارائه شده در مقالات مقایسه شد. رفتار فازها طی فرآیند انتقال حرارت در زوایای مختلف متفاوت بود. مشاهده شد که عدد ناسلت، اختلاف دمای هوا و اختلاف دمای ذرات در نقاط مشخص شده در لوله با زاویه ۴۵ درجه نسبت به دیگر لوله‌ها بیشتر است. افزایش در نرخ ورود ذرات جامد منجر به افزایش در اختلاف دمای هوا شد. از طرفی افزایش نرخ ورود ذرات به ترتیب منجر به افزایش و کاهش اختلاف دمای ذرات در نرخ ورودی پایین و بالا شد.

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